

# Lecture 2: Nonlinear Phenomena

## Review

Pendulum example:

"assume force resisting motion of the pendulum, proportional to the speed of the mass"

$$F_{ext} = -kV$$

$$= -kl\dot{\theta}$$

$$\frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} - \frac{\partial \mathcal{L}}{\partial \theta} = \tau_{ext}$$

$$ml^2\ddot{\theta} + mgl\sin\theta = -kl^2\dot{\theta}$$

$$\ddot{\theta} = -\frac{k}{m}\dot{\theta} - \frac{g}{l}\sin\theta$$



$$x = \begin{bmatrix} \theta \\ \dot{\theta} \end{bmatrix} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}, \quad \dot{x} = \begin{bmatrix} \dot{\theta} \\ \ddot{\theta} \end{bmatrix} = \begin{bmatrix} x_2 \\ -\frac{k}{m}x_2 - \frac{g}{l}\sin x_1 \end{bmatrix}$$

Steps:

① What are the equilibrium points?

$$x^* = f(x^*) = 0$$

$$x^* = \left\{ \underbrace{(0, 0)}_{x_1^*, x_2^*}, \underbrace{(\pi, 0)}_{x_1^*, x_2^*} \right\}$$

② Linearize our system about  $x^*$ .

$$\dot{x} = Ax$$

$$\frac{\partial f}{\partial x} \Big|_{x^*}$$

$$\frac{\partial f}{\partial x} = \begin{bmatrix} 0 & 1 \\ -\frac{g}{l}\cos x_1 & -\frac{k}{m} \end{bmatrix}$$

$$\frac{\partial f}{\partial x} \Big|_{x^* = (0, 0)} = \begin{bmatrix} 0 & 1 \\ -\frac{g}{l} & -\frac{k}{m} \end{bmatrix}$$

$\Rightarrow \text{Re}(\lambda_i(A)) < 0$   
 $\Rightarrow$  stable!

$$\frac{\partial f}{\partial x} \Big|_{x^* = (\pi, 0)} = \begin{bmatrix} 0 & 1 \\ \frac{g}{l} & -\frac{k}{m} \end{bmatrix}$$

$\Rightarrow \text{Re}(\lambda_1(A)) < 0$   
 $\text{Re}(\lambda_2(A)) > 0 \Rightarrow$  unstable

Recall Jacobian

$$\frac{\partial f}{\partial x} := \begin{bmatrix} \frac{\partial f_1}{\partial x_1} & \frac{\partial f_1}{\partial x_2} \\ \frac{\partial f_2}{\partial x_1} & \frac{\partial f_2}{\partial x_2} \end{bmatrix}$$

Recall

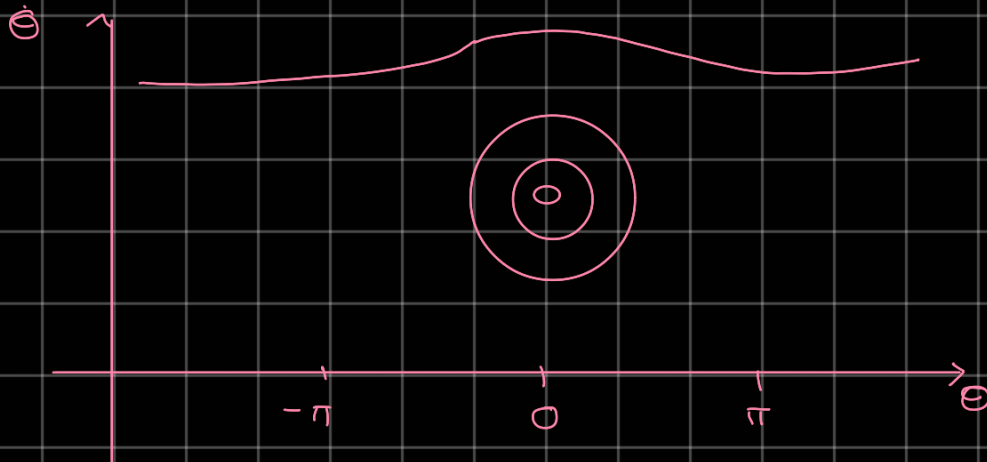
- $\rightarrow$  negative eigenvalues imply stability
- $\rightarrow$  positive eigenvalues imply instability

Takeaway: Linearization is a valid approach for analyzing the behavior of the nonlinear systems.

But the limitations are that

I. only valid locally

II. nonlinear dynamics are richer



### Essential Nonlinear Phenomena

#### ① Finite Escape Time

(the state  $x$  goes to  $\infty$  in  $\infty$  time)

Example  $\dot{x} = x^2$

↳ For a linear system, this would never happen

$$\dot{x} = \frac{dx}{dt} = x^2$$

$$\frac{1}{x^2} dx = dt$$

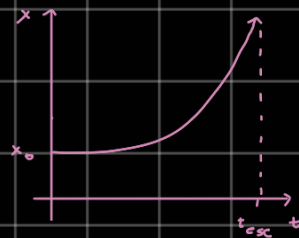
$$\int \frac{1}{x^2} dx = \int dt$$

$$-\frac{1}{x} = t + C$$

$$x(t) = \frac{1}{C-t}, \quad x(0) = \frac{1}{C}$$

$$x(t) = \frac{1}{\frac{1}{x_0} - t}$$

$$t_{\text{escape}} = \frac{1}{x_0}$$



#### ② Multiple Isolated Equilibria

Linear system:

- either unique  $x^*$  or continuum

Pendulum: two isolated  $x^*$

"multi-stable" systems

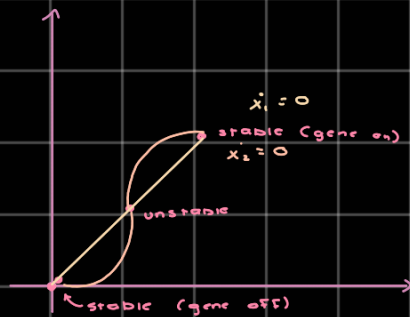
Example: Bi-stable switching

$$\dot{x}_1 = -ax_1 + x_2 \quad \leftarrow \text{concentration of protein}$$

$$\dot{x}_2 = \frac{x_1^2}{1+x_1^2} - bx_2 \quad \leftarrow \text{mRNA}$$

$$a > 0, b > 0$$

→ gene stimulation by transcription  $\frac{x_1^2}{1+x_1^2}$



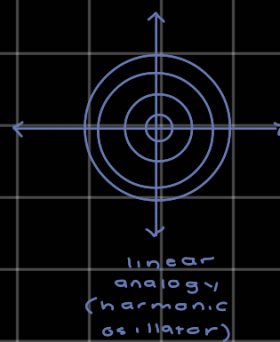
③ Limit cycles → "like a trajectory of equilibrium points, in state space"

- linear oscillators → continuum of periodic orbits

every "circle"  $\dot{x} = Ax$   $A = \begin{bmatrix} 0 & -\beta \\ \beta & 0 \end{bmatrix}$



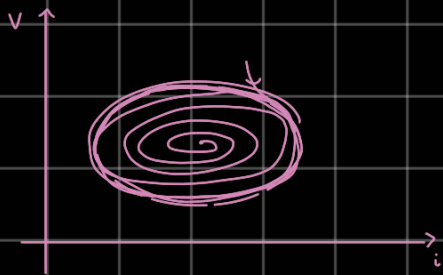
Nonlinear  
↳ idea



linear  
analogy  
(harmonic  
oscillator)

Example: van der Pol oscillator

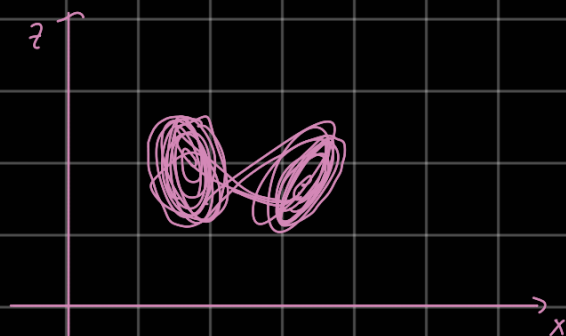
→ system w/ self-sustained oscillations



④ Chaos: irregular oscillations

(never exactly repeating)

Example: Lorenz System



→ For a continuous-time, time-invariant system

↳ need  $n \geq 3$  chaos

$n=2$  Poincaré - Bendixson Theorem: guarantees regular behavior  
 → only applies to time-invariant systems

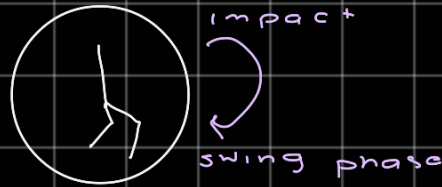
→ for time-varying, only need  $n \geq 2$

→ For discrete time

$\delta x_k, n \geq 1$  is sufficient for chaos  
 (bc of lack of continuity assumptions)

⑤ Multiple Modes of Behavior (Hybrid systems)

Ex: walking!



$$\dot{x} = f(x) + g(x)u$$

$$H = \begin{cases} \dot{x} = f(x) + g(x)u & x \notin S \\ x^+ = \Delta(x) & x \in S \end{cases}$$

Planar (second-order) Dynamical systems

$$x = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}$$

$$\dot{x}_1 = F_1(x_1, x_2)$$

$$\dot{x}_2 = F_2(x_1, x_2)$$

• solution trajectories are "curves" in the plane

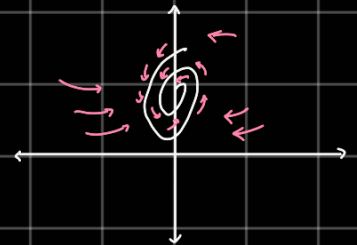
$$x(t) = (x_1(t), x_2(t))$$

↳ either a phase plane or a state plane

• Family of all trajectories (solution curves)

is called the **phase portrait**

↳  $f(x)$  is the vector field on the plane



→ Let's look at phase portraits for linear systems!

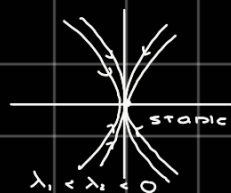
$$\dot{x} = Ax$$

Jordan Form:  $J = T^{-1}AT$

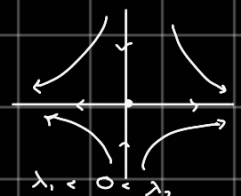
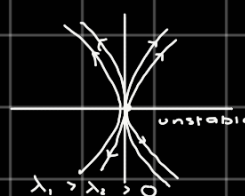
change of coordinates:  $z = T^{-1}x$

3 options for the Jordan form:

①  $\begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \quad D$



$$\begin{aligned} \dot{z}_1 &= \lambda_1 z_1 \\ \dot{z}_2 &= \lambda_2 z_2 \end{aligned}$$



$$\begin{bmatrix} \lambda & \kappa \\ 0 & \lambda \end{bmatrix}, \begin{bmatrix} \kappa & -\beta \\ \beta & \kappa \end{bmatrix}$$